

## Dependence of 3D Self-correlation Level Contours on the Scales in the Inertial Range of Solar Wind Turbulence

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## Abstract

The self-correlation level contours at the  $10^{10}$  cm scale reveal a 3D isotropic feature in the slow solar wind and a quasi-anisotropic feature in the fast solar wind. However, the  $10^{10}$  cm scale is approximately near the low-frequency break (outer scale of turbulence cascade), especially in the fast wind. How the self-correlation level contours behave with dependence on the scales in the inertial range of solar wind turbulence remains unknown. Here we present the 3D self-correlation function level contours and their dependence on the scales in the inertial range for the first time. We use data at 1 au from instruments on the *Wind* spacecraft in the period 2005–2018. We show the 3D isotropic self-correlation level contours of the magnetic field in the inertial range of both slow and fast solar wind turbulence. We also find that the self-correlation level contours of the velocity in the inertial range present 2D anisotropy with an elongation in the perpendicular direction and 2D isotropy in the plane perpendicular to the mean magnetic field. These results indicate differences between the magnetic field and the velocity, providing new clues to interpret the solar wind turbulence on the inertial scale.

*Unified Astronomy Thesaurus concepts:* Solar wind (1534); Interplanetary turbulence (830); Magnetic fields (944); Space plasmas (1544)

## 1. Introduction

The magnetic field and the velocity both display broadband fluctuations in the solar wind. The ubiquitous observation of Kolmogorov-like magnetic and velocity power spectra suggests the existence of a turbulent energy cascade in the inertial range (Frisch 1995; Tu & Marsch 1995; Bruno & Carbone 2013). At the large-scale side of the inertial range, the fast solar wind often presents a robust 1/f scaling injection range (Horbury et al. 1996; Matthaeus et al. 2007; Bruno & Carbone 2013) for both the magnetic field and the velocity. The low-frequency break between the inertial range and the injection range is around  $10^{-3}$  Hz, a typical value for the fast wind at 1 au (Bruno & Carbone 2013; Bruno et al. 2019). In the slow solar wind, the 1/f scaling is also present for the magnetic field with a smaller low-frequency break around  $10^{-4}$  Hz. However, the velocity spectrum keeps the Kolmogorov-like scaling throughout the analyzed frequency range by Bruno et al. (2019).

Matthaeus et al. (1990) developed the 2D self-correlation function method to study the solar wind turbulence and obtained the famous "Maltese cross." Dasso et al. (2005) applied the same method and found that the anisotropy behaves differently for the slow wind and the fast wind shown by the self-correlation function level contours analyzing the two-day intervals measured by the Advanced Composition Explorer (ACE) spacecraft at 1 au. A 2D self-correlation function is also constructed by analyzing the simultaneous measurements from the Cluster-4 spacecraft, showing anisotropic characteristics at small scales close to ion kinetic scales for both solar wind and magnetosheath turbulences (Osman & Horbury 2006; He et al. 2011). Wang et al. (2019) extended this study using the same data set and found that the anisotropy disappears and becomes 2D isotropic for both the slow wind and the fast wind, with the interval durations decreasing from 2 days, to 1 day, to 10 hr, to 2 hr, to 1 hr. Wu et al. (2019) further extended the selfcorrelation function level contours analysis using 1 hr intervals

observed by *WIND* spacecraft. They show 3D isotropic selfcorrelation function level contours in the slow wind and 3D quasi-isotropic self-correlation function level contours in the fast wind. The contour scale of 1 hr intervals is around  $10^{10}$  cm near the low-frequency break scale. However, the feature of the self-correlation function level contours and its dependence on the scales in the inertial range remains unknown.

In the present study, we perform a 3D self-correlation function level contour analysis on the *WIND* spacecraft measurements using intervals with durations = 1 hr, 30 minutes, and 10 minutes. We briefly introduce the method in Section 2 and present our observational results in Section 3. In Section 4, we draw our conclusions.

### 2. Data and Method

We briefly describe the method used in the analysis; more details can be found in Wu et al. (2019). We use the magnetic field data with a cadence of  $\Delta = 3$  s from the magnetic field investigation (Lepping et al. 1995) and the plasma data with the same time resolution (3 s) from the three-dimensional plasma analyzer (Lin et al. 1995) on board the *WIND* spacecraft in the period 2005–2018. The data set was cut into intervals with a duration of *T*, where T = 1 hr, 30 minutes, and 10 minutes, respectively. These intervals were conserved for further investigation, and contain a less than 5% data gap and max[ $|\delta B_j| < 2$  nT, max[ $|\delta V_j| < 20$  km, where *j* indicates *x*, *y*, *z* components in the geocentric-solar-ecliptic coordinate system, and  $\delta$  means the variation between every 3 s.

The two-time-point self-correlation function for each interval *i* is defined as

$$R_U(i,\tau) = \langle \delta \boldsymbol{U}(t) \cdot \delta \boldsymbol{U}(t+\tau) \rangle, \qquad (1)$$

where  $\tau = 0, \Delta, 2\Delta, ..., T/2$  is the time lag,  $\langle \rangle$  denotes an ensemble time average, and  $\delta U$  is the time series removing a linear trend for either magnetic field **B** or velocity **V**. We



Figure 1. (a): averaged normalized self-correlation functions  $R_{bb}(r)$  of 30 minute magnetic field data with the standard error bars of  $r_{level}$  for a given  $R_{bb}$ . The solid (dashed) lines are for the slow (fast) wind. Red, blue, and yellow indicate the  $r_{\parallel}$ ,  $r_{\perp 1}$ , and  $r_{\perp 2}$  directions, respectively. (b): averaged normalized self-correlation functions  $R_{vv}(r)$  of 30 minute velocity data, in the same manner as in (a). (c): same as (a) for 10 minute magnetic field data. (d): same as (b) for 10 minute velocity data.

normalize the self-correlation function using the zero time-lag self-correlation  $R_{uu}(i, \tau) = R_U(i, \tau)/R(i, 0)$ . We further obtain the spatial lag *r* using the Taylor hypothesis (Taylor 1938)  $r = \tau V_{SW}$ , where  $V_{SW}$  is the mean flow velocity in the corresponding interval *i*.

The 3D coordinate system is constructed using the mean magnetic field  $B_0$  and the maximum variance direction L obtained by the minimum-variance analysis (MVA) method (Sonnerup & Cahill 1967). The  $r_{\parallel}$  and  $r_{\perp 2}$  components are defined as the mean magnetic field  $B_0$  and the projection of L in

the plane perpendicular to  $B_0$ .  $r_{\perp 1} = r_{\parallel} \times r_{\perp 2}$ . We calculate the angle  $\theta_{\rm VB}$  between  $V_{\rm SW}$  and  $B_0$  and the angle  $\phi_{\rm L}$  between the  $r_{\perp 2}$  direction and the component of  $V_{\rm SW}$  perpendicular to  $B_0$ for each interval *i*.

We divide these intervals into the slow wind  $(V_{\rm SW} < 400 \,\rm km \, s^{-1})$  and the fast wind  $(V_{\rm SW} > 500 \,\rm km \, s^{-1})$  and study their 3D self-correlation level contours separately. For T = 30 minutes, we obtain 55331 intervals in the slow wind and 10733 intervals in the fast wind. For T = 10 minutes, the numbers are 217830 and 63656.  $\theta_{\rm VB}$  and  $\phi_{\rm L}$  are binned into



Figure 2. The 3D self-correlation level contour surface at level  $R_{uu} = 0.368$  of the 10 minute data for: (a) the magnetic field in the slow wind; (b) the magnetic field in the fast wind; (c) velocity in the slow wind; and (d) the velocity field in the fast wind. The color represents the distances from the origin  $r_{level}$  [10<sup>10</sup> cm]. The dashed red (blue) lines in  $r_{\perp 1} = -0.70$  plane are projections of the intersection lines of the surface with two planes  $r_{\perp 1} = A1$  (A2), where A1 and A2 are shown in the legends with the corresponding colors in the corresponding panel, and the same for the other two planes.

 $15^\circ$  bins and the average of the normalized self-correlation functions is calculated as

$$R_{uu}(\theta_{VB}^{m}, \phi_{L}^{n}, r) = \frac{1}{n(\theta_{VB}^{m}, \phi_{L}^{n})} \sum_{\substack{\theta_{VB}^{n} = 7.5 < =\theta_{VB}(i) < \theta_{VB}^{m} + 7.5, \\ \phi_{L}^{n} - 7.5 < =\phi_{1}(i) < \phi_{L}^{n} + 7.5}} R_{uu}(i, r), \qquad (2)$$

where  $n(\theta_{VB}^{m}, \phi_{L}^{n})$  denotes the number of intervals in the corresponding bin, and  $\theta_{VB}^{m} = 15^{\circ}m + 7^{\circ}.5$ ;  $\phi_{L}^{n} = 15^{\circ}n + 7^{\circ}.5$ ; m, n = 0, 1, 2, ..., 5.

We obtain 36 averaged self-correlation functions for 36  $(\theta_{\rm VB}, \phi_{\rm L}) = 15^{\circ} \times 15^{\circ}$  bins. Figure 1 shows the averaged self-correlation functions for T = 30 minutes and T = 10 minutes in the  $r_{\perp 1}$  (75° < =  $\theta_{\rm VB}$  < = 90°, 75° < =  $\phi_{\rm L}$  < = 90°),  $r_{\perp 2}$  (75° < =  $\theta_{\rm VB}$  < = 90°, 0° < =  $\phi_{\rm L}$  < 15°), and  $r_{\parallel}$  (0° < =  $\theta_{\rm VB}$  < 15°, 0° < =  $\phi_{\rm L}$  < = 90°) directions.

In the left panels of Figure 1, we present the averaged magnetic self-correlation functions with standard error bars in both the slow wind (solid lines) and the fast wind (dashed lines). The functions of the three directions are almost overlapped with each other for both 30 minute and 10 minute intervals, especially for the slow wind, indicating isotropy of the self-correlation functions. In the right panels of Figure 1, we show the averaged self-correlation functions of the velocity with standard error bars for both the slow wind and the fast wind. The velocity functions of the three directions have more difference with each other than the magnetic field. In general, the parallel function is smaller than both of the perpendicular functions and the perpendicular functions behave similar with each other. For both the magnetic field and the velocity, the correlation functions of the slow wind decrease more rapidly than that of the fast wind for both 30 minute and 10 minute intervals. This difference between the slow and fast wind has already been shown for 1 hr intervals in Wu et al. (2019),



Figure 3. The 3D self-correlation level contour surface at level  $R_{uu} = 0.368$  of the 10 minute data with the same panel and line styles as Figure 2, except the projection planes are  $r_{\parallel} = -0.30$ ,  $r_{\perp 1} = -0.30$ , and  $r_{\perp 2} = -0.30$ .

where 23083 intervals in the slow wind and 3347 intervals in the fast wind are investigated.

For each bin  $(\theta_{\rm VB}, \phi_{\rm L})$ , we calculate *r* at level  $R_{uu}(\theta_{\rm VB}, \phi_{\rm L}, r) = 1/e \approx 0.368$  and denote the result as  $r_{\rm level}$ . We transform  $(\theta_{\rm VB}, \phi_{\rm L}, r_{\rm level})$  into  $(r_{\perp 1}, r_{\perp 2}, r_{\parallel})$  as

$$r_{\perp 1} = r_{\text{level}} \sin \theta_{\text{VB}} \sin \phi_{\text{L}},\tag{3}$$

$$r_{\perp 2} = r_{\text{level}} \sin \theta_{\text{VB}} \cos \phi_{\text{L}},\tag{4}$$

$$r_{\parallel} = r_{\text{level}} \cos \theta_{\text{VB}}.$$
 (5)

We also obtain  $r_{\text{level}}(i)$  for each interval *i* at level  $R_{uu}(i, r) = 1/e \approx 0.368$ . We define two ratios  $r_{\parallel}^c/r_{\perp}^c$  and  $r_{\perp 2}^c/r_{\perp 1}^c$ , where

$$r_{\parallel}^{c} = \frac{1}{n(r_{\parallel}^{c})} \sum_{\substack{0 < =\theta_{\rm VB}(i) < 15, \\ 0 < =\phi_{\rm L}(i) < 90}} r_{\rm level}(i), \tag{6}$$

$$r_{\perp}^{c} = \frac{1}{n(r_{\perp}^{c})} \sum_{\substack{75 < =\theta_{\rm VB}(i) < 90, \\ 0 < =\phi_{\rm L}(i) < 90}} r_{\rm level}(i), \tag{7}$$

$$r_{\perp 2}^{c} = \frac{1}{n(r_{\perp 2}^{c})} \sum_{\substack{60 < =\theta_{\rm VB}(i) < 90, \\ 0 < =\phi_{\rm L}(i) < 15}} r_{\rm level}(i), \tag{8}$$

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and

$$r_{\perp 1}^{c} = \frac{1}{n(r_{\perp 1}^{c})} \sum_{\substack{60 < =\theta_{\rm VB}(i) < 90, \\ 75 < =\phi_{\rm L}(i) < 90}} r_{\rm level}(i).$$
(9)

The ratio  $r_{\parallel}^c/r_{\perp}^c$  describes the  $r_{\text{level}}$  difference between the parallel and the perpendicular direction, and the ratio  $r_{\perp 2}^c/r_{\perp 1}^c$  describes the anisotropy in the perpendicular plane. The result is shown in the next section.

#### 3. Results

Figure 2 shows 3D self-correlation level contour surfaces at level  $R_{uu} = 0.368$  for 30 minute intervals. The spatial lag scale for  $R_{uu} = 0.368$  is around 10<sup>9</sup> cm for 30 minutes, which is on the inertial scale of solar wind turbulence at 1 au. In



**Figure 4.** Left panel:  $r_{\parallel}^c/r_{\perp}^c$  for intervals with different durations. The red (black) lines are for the slow (fast) wind. The solid (dashed) lines indicate the magnetic field (velocity) results. From left to right, the vertical dashed lines correspond to 2 days, 1 day, 10 hr, 2 hr, 1 hr, 30 minutes, and 10 minutes, respectively. The ratios shown as hollow circles and hollow triangles are from Wang et al. (2019) with *ACE* measurements at level = 0.8. While the solid circles and solid triangles are from our *WIND* observations at level  $R_{uu} = 0.368$ . Right panel:  $r_{\perp 2}^c/r_{\perp 1}^c$  for intervals with different durations at level  $R_{uu} = 0.368$  with the same line styles as in the left panel. From left to right, the vertical dashed lines correspond to 1 hr, 30 minutes, and 10 minutes, respectively.

Figure 2(a), the slow wind magnetic field self-correlation level contour surface is spherical. The projection closed curves on the 2D plane are plotted to help visualize the 3D feature. In Figure 2(b), the fast wind magnetic field self-correlation level contour surface shows a weak elongation along  $r_{\perp 2}$ . In Figure 2(c), the slow wind velocity self-correlation level contour surface presents a weak elongation in the perpendicular plane. In Figure 2(d), the fast wind velocity field selfcorrelation level contour surface has a similar shape with that of magnetic field. The size difference of the surface remains for 30 minute intervals, as seen for the 1 hr intervals shown in Wu et al. (2019): the  $r_{\text{level}}$  is longer for the fast wind than for the slow wind and longer for the magnetic field than for the velocity. We also analyzed the 3D self-correlation level contours in the LMN coordinate system, constructed by MVA analysis: L is the same directions of maximum variance as those used for the construction of the 3D coordinate system described in Section 2; N is the directions of minimum variance, and M is directions of immediate variance. The 3D features of the self-correlation level contours in the LMN coordinate system (not shown) are the same as those shown in Figure 2.

Figure 3 shows 3D self-correlation level contour surfaces at level  $R_{uu} = 0.368$  for 10 minute intervals. The size of surface reaches  $6 \times 10^8$  cm in the slow wind for the velocity. The magnetic field self-correlation level contour surfaces for both the slow wind and the fast wind are almost spherical, suggesting isotropy for the magnetic field, as shown in Figures 3(a) and (b). The elongation in the perpendicular plane of the velocity field self-correlation function contour surfaces grows in both the slow wind and the fast wind seen in Figures 3(c) and (d).

In the left panel of Figure 4, we show  $r_{\parallel}^c/r_{\perp}^c$  with interval durations = 1 hr, 30 minutes, and 10 minutes for the magnetic field (solid circles) and the velocity field (solid triangles) measured by the *Wind* spacecraft for both the slow (red) and fast (black) winds. For comparison, the results of *ACE* 

observations from Wang et al. (2019) at level  $R_{uu} = 0.8$  with time durations = 2 days, 1 day, 10 hr, 2 hr, 1 hr are shown here as hollow circles and hollow triangles. The results of 1 hr intervals in our work are not exactly the same as those in Wang et al. (2019). That may attribute to the different data sets. The  $r_{\parallel}^{c}/r_{\parallel}^{c}$  of the magnetic field in both the slow wind and the fast wind with interval durations = 1 hr, 30 minutes, and 10 minutes are all very close to 1, indicating 2D isotropy, as shown by the self-correlation level contour on the inertial scale. The ratios of the velocity field are less than 1, especially in the slow wind. They are 0.72 and 0.86 for 10 minute intervals in the slow wind and fast wind obtained by averaging  $r_{level}$ . We calculate the ratios directly using the  $r_{level}$  calculated by the correlation functions shown in Figure 1 at  $R_{\nu\nu} = 0.368$  and they are 0.73 and 0.86. These results suggest an elongation along the perpendicular direction for the velocity on the inertial scale. The right panel of Figure 4 shows  $r_{\perp 2}^c/r_{\perp 1}^c$ , suggesting isotropy in the perpendicular plane for both the magnetic field and the velocity field in the slow wind. In the fast wind, the elongations along  $r_{\perp 2}$  for 1 hr disappear for 10 minutes, indicating isotropy in the perpendicular plane deep into the inertial scale.

#### 4. Discussion and Conclusions

We present 3D self-correlation level contours and their dependence on the timescale in the inertial range using *WIND* measurements at 1 au during 13 years from 2005 to 2018. We analyze the self-correlation level contours of the magnetic field and the velocity in both the slow wind and the fast wind for intervals with durations = 30 minutes and 10 minutes. We use two ratios  $r_{\perp 2}^{c}/r_{\perp 1}^{c}$  to describe the anisotropy and show their dependence on the timescale of the interval duration. The 3D self-correlation level contours of the magnetic field present isotropy in both the slow solar wind and the fast solar wind for both the 30 minute intervals and the 10 minute intervals, which corresponds to approximately a scale of 10<sup>9</sup>

cm on the inertial scale. However, the 3D self-correlation level contours of the velocity indicate an elongation in the direction perpendicular to the mean magnetic field, and 2D isotropy in the plane perpendicular to the mean magnetic field. The behaviors of the magnetic field and the velocity and their differences are new.

Carbone et al. (1995) developed a model for 3D magnetic field correlation spectra and reconstructed the shape of the selfcorrelation level contours of "Maltese cross" using the minimum-variance framework. Our new results are consistent with the "Maltese cross" and this inconsistency requires further study. It is also hard to understand our results under the framework of critical balance theory (Goldreich & Sridhar 1995; Boldyrev 2006), which predicts strong anisotropy of the MHD turbulence spectrum that is inconsistent with some local observations of structure functions in the solar wind (Chen et al. 2012; Verdini et al. 2018). The existing theories cannot explain our results and we cannot provide an exhaustive explanation with our current level of understanding and numerical simulations. These results open a new window into interpreting solar wind turbulence on the inertial scale.

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#### References

- Boldyrev, S. 2006, PhRvL, 96, 115002
- Bruno, R., & Carbone, V. 2013, LRSP, 10, 2
- Bruno, R., Telloni, D., Sorriso-Valvo, L., et al. 2019, A&A, 627, A96
- Carbone, V., Malara, F., & Veltri, P. 1995, JGR, 100, 1763
- Chen, C. H. K., Mallet, A., Schekochihin, A. A., et al. 2012, ApJ, 758, 120
- Dasso, S., Milano, L. J., Matthaeus, W. H., & Smith, C. W. 2005, ApJL, 635 L181
- Frisch, U. 1995, Turbulence. The legacy of A.N. Kolmogorov (Cambridge: Cambridge Univ. Press)
- Goldreich, P., & Sridhar, S. 1995, ApJ, 438, 763
- He, J.-S., Marsch, E., Tu, C.-Y., et al. 2011, JGRA, 116, A06207
- Horbury, T. S., Balogh, A., Forsyth, R. J., & Smith, E. J. 1996, A&A, 316, 333
- Lepping, R. P., Acũna, M. H., Burlaga, L. F., et al. 1995, SSRv, 71, 207
- Lin, R. P., Anderson, K. A., Ashford, S., et al. 1995, SSRv, 71, 125
- Matthaeus, W. H., Breech, B., Dmitruk, P., et al. 2007, ApJL, 657, L121
- Matthaeus, W. H., Goldstein, M. L., & Roberts, D. A. 1990, JGR, 95, 20673 Osman, K. T., & Horbury, T. S. 2006, ApJL, 654, L103
- Sonnerup, B. U. O., & Cahill, L. J., Jr. 1967, JGR, 72, 171
- Taylor, G. I. 1938, RSPSA, 164, 476
- Tu, C.-Y., & Marsch, E. 1995, SSRv, 73, 1
- Verdini, A., Grappin, R., Alexandrova, O., & Lion, S. 2018, ApJ, 853, 85
- Wang, X., Tu, C., & He, J. 2019, ApJ, 871, 93
- Wu, H., Tu, C., Wang, X., He, J., & Wang, L. 2019, ApJ, 882, 21



# Erratum: "Dependence of 3D Self-correlation Level Contours on the Scales in the Inertial Range of Solar Wind Turbulence" (2019, ApJL, 883, L9)

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In the caption of Figure 2 in the published article, the time "30 minute" in the first sentence is mistakenly written as "10 minute." A corrected caption of Figure 2 has been provided here.

In the caption of Figure 4 of the published article, the last sentence is "From left to right, the vertical dashed lines correspond to 1 hr, 30 minutes, and 10 minutes, respectively." However, there are no vertical dashed lines. In the corrected caption of Figure 4 provided here, this sentence is rewritten as "From left to right, the timescales correspond to 1 hr, 30 minutes, and 10 minutes, respectively."





**Figure 2.** The 3D self-correlation level contour surface at level  $R_{uu} = 0.368$  of the 30 minute long data for the (a) magnetic field in the slow wind; (b) magnetic field in the fast wind; (c) velocity in the slow wind; and (d) velocity in the fast wind. The color represents the distances from the origin  $r_{\text{level}} [10^{10} \text{ cm}]$ . The dashed red (blue) lines in  $r_{\perp 1} = -0.70$  plane are projections of the intersection lines of the surface with two planes  $r_{\perp 1} = A1$  (A2), where A1 and A2 are shown in the legends with the corresponding colors in the corresponding panel, same for the other two planes.



Figure 4. Left panel:  $r_{\parallel}^c/r_{\perp}^c$  for intervals with different duration. The red (black) lines are for the slow (fast) wind. The solid (dashed) lines indicate the magnetic field (velocity) results. From left to right, the timescales correspond to 2 days, 1 day, 10 hr, 2 hr, 1 hr, 30 minutes, and 10 minutes, respectively. The ratios shown in hollow circles and hollow triangles are from Wang et al. (2019) with ACE measurements at level = 0.8. While the solid circles and solid triangles are from our WIND observations at level  $R_{uu} = 0.368$ . Right panel:  $r_{\perp 2}^c/r_{\perp 1}^c$  for intervals with different duration at level  $R_{uu} = 0.368$  with the same line styles as in the left panel. From left to right, the timescales correspond to 1 hr, 30 minutes, respectively.

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#### References

Wang, X., Tu, C., & He, J. 2019, ApJ, 871, 93